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Chordwise Divergence of Rectangular Flat Plate Wings at Supersonic Speeds

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Nomenclature

speed of sound aeroelastic lift efficiency wing chord generalized coordinates Young's modulus spanwise mode shapes $f(\vartheta), F(\vartheta)$ $q_{(\eta)}, G_{(\eta)}$ chordwise mode shapes wing span (root to tip) L_i lift/unit incidence of element j integers m, nMMach number generalized coordinates r_m wing thickness Ŵ wing deflection incidence of rigid wing α change in incidence, due to deformation, of element j α_{j} $(M^2-1)^{1/2}$ β nondimensional chordwise coordinate η nondimensional spanwise coordinate air density deflection function

Introduction

IN most of the published papers dealing with the chordwise divergence of rectangular flat plate wings clamped along their root chord it has been assumed that the spanwise mode shape is relatively unimportant, compared with the chordwise mode shape, and can be fixed arbitrarily. The wide variation in the results obtained (Fig. 1) has shown the need for analyses not embodying this assumption. Such analyses have been made by Rhodes1 using completely general assumed deflection modes and Ackeret aerodynamic theory.

Previous studies using Ackeret aerodynamic theory have been made by Biot,^{2,3} Broadbent⁴ and Martin.⁵ For a twodimensional wing with no spanwise bending, Biot² derived an exact stability criterion for chordwise divergence which may be written

$$(E/\rho a^2)(t/l)^3 > 0.43(C/l)^3 M^2/\beta$$
 (1)

The stabilizing influence of the anticlastic curvature induced by spanwise bending was shown in later papers by

 $\varphi(\zeta,\eta)$

Biot,^{2,3} and the influence of the assumed spanwise mode on the results was also indicated. Thus, if

$$W = f_{(\zeta)} \cdot g_{(\eta)} = \zeta^2 (3 - \zeta) \sum_{n=0}^{6} d_n \eta^n$$
 (2)

divergence was not possible for flat plates of aspect ratio 3

More recent studies by the authors using as modes

$$W = \zeta^{2}(3 - \zeta) \sum_{n=0}^{2} d_{n} \eta^{n} + \zeta(2 - \zeta) \sum_{m=1}^{2} r_{m} \eta^{m}$$
i. e.
$$W = f_{(\zeta)} \cdot g_{(\eta)} + F_{(\zeta)} G_{(\eta)}$$
(3)

showed that divergence was possible at least for aspect ratios up to 2 (Fig. 1). Note that function $F \cdot G$ only satisfies the condition $\partial W/\partial \zeta = 0$ at $\eta = 0$.

Also shown in Fig. 1 are results using the modes

$$W = d_1 \zeta^3 \eta^2 + d_2 \zeta^2 \eta \tag{4}$$

as well as those from Refs. 1-5.

Broadbent⁴ assumed that the plate deformed in a cylindrical mode about a generator through the root leading edge swept at 45° to the freestream. He obtained the following stability criterion:

$$(E/\rho a^2)(t/l)^3 > 0.3M^2/\beta$$
 (5)

Martin⁵ assumed

$$W = (d_1 + d_2)\zeta^3\eta^2 + d_3[(\zeta^2/2) - (\zeta^3/3) + (\zeta^4/12)] + d_4\eta(2\zeta - \zeta^2)$$
 (6)

where d_1 and d_2 are generalized coordinates for chordwise deformation forward of, and aft of, the mid chord, respectively. The results indicated that divergence was possible for low aspect ratio plate wings at high supersonic speeds.

Hancock⁶ used an approximate, linearized, supersonic pressure distribution, which is much more general than Ackeret theory, and compared his results with Ref. 7, which used linearized slender wing theory and is only valid for low aspect ratios at $M \simeq 1$. Hancock showed, using as assumed modes

$$W = \frac{1}{2.6} (45\zeta^2 - 20\zeta^3 + \zeta^6) g_{(\eta)} \tag{7}$$

that chordwise divergence was not possible outside of the

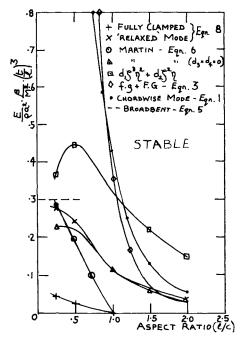


Fig. 1 Flat plate stability boundaries for $M \gg 1$.

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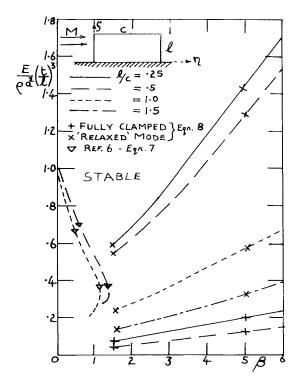


Fig. 2 Flat plate stability boundaries for $M \ge 1$ ($\beta \ge 0$).

range 1 < M < 1.8 (Fig. 2). Also, differences in the stability criteria obtained in Refs. 6 and 7 for M = 1 could be explained completely by the differences in the assumed spanwise mode shapes.

Solutions Using General Deflection Modes

Rhodes¹ assumed that

$$W = C\varphi_{(t,n)} \tag{8}$$

where

$$\varphi_{(\zeta,\eta)} = \sum_{m=0}^{4} \sum_{n=0}^{4} C_{mn} \zeta^{m} \eta^{n}$$

The clamped root boundary conditions at $\zeta = 0$, for all η , imply that values of m = 0, 1 are excluded from Eq. (8). A stability determinant of order 15 is obtained, and the results for the stability criteria obtained are given in Figs. 1 and 2. Divergence was not possible at high supersonic speeds for aspect ratios greater than unity, and even for lower aspect ratios the critical plate thickness is so low as to virtually preclude divergence.

Results were also obtained for a relaxed mode in which the root boundary condition $\partial W/\partial \zeta = 0$ was only enforced at $\eta = 0$, the root leading edge. A stability determinant of order 19 was obtained, and the results are plotted in Figs. 1 and 2. The relaxed modes give larger areas of instability than the fully clamped modes and would suggest that divergence was possible at high supersonic speeds for aspect ratios at least up to 2.

Thus, as indicated by the results using Eqs. (2) and (3), relaxation of the spanwise mode shape through the root boundary conditions can indicate larger areas of instability than more exact analyses would give.

Solutions Using Structural Influence Coefficients

In Ref. 6 experimental influence coefficients have been obtained for a plate of aspect ratio 0.5 and used in a direct matrix analysis for divergence. The results confirm those obtained by energy analyses using assumed modes, Eq. (7), for values of $\beta = 0.45$ and 1.0 in Fig. 2.

Both experimental and theoretical structural influence coefficients have been used in Ref. 1, for plates of aspect ratio 2 and 1, respectively, with Ackeret aerodynamic theory.

The aeroelastic problem considered requires the solution of

$$\alpha_i = \sum_{j=1}^{N} L_j a_{ij} (\alpha + \alpha_j)$$
 (9)

and the aeroclastic lift efficiency A_E (i.e., the ratio of lift on the elastic wing to that on the rigid wing)

$$A_E = \sum_{j=1}^{N} L_j(\alpha + \alpha_j) / \sum_{j=1}^{N} L_j \alpha$$
 (10)

becomes, if $L_i = L = \text{const}$,

$$A_E = 1 + \frac{1}{N} \sum_{j=1}^{N} \frac{\alpha_j}{\alpha}$$
 (11)

The calculations indicated that for aspect ratios 1 and 2 (with N=25 and 50, respectively) the aeroclastic lift efficiency A_E is always less than unity and divergence (when $A_E \rightarrow \alpha$) is not possible. These results confirm those obtained with the fully clamped mode in the energy analyses described earlier.

Conclusions

The foregoing studies confirm that divergence is most unlikely at high supersonic speeds, for rectangular flat plate wings, on the basis of linear structural theory and Ackeret aerodynamic theory. Also, the choice of assumed modes in divergence analyses is very critical and completely general spanwise and chordwise mode shapes should be considered which satisfy all of the root boundary conditions.

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